

1. Lathi 1.1-9

This  $x(t)$  is not a periodic signal.

Let's compute the energy first...

$$E_x = \int_{-\infty}^{\infty} x^2(t) dt = \int_0^{\infty} x^2(t) dt \quad \text{because } x(t) = 0 \text{ for all } t < 0$$

$$= \int_0^2 x^2(t) dt + \int_2^5 x^2(t) dt + \int_5^9 x^2(t) dt + \dots$$

In each of these integration regions,  $x(t)$  is on for the first second and off for the remaining time. Hence, each integral evaluates to be equal to one.

$$E_x = \sum_{n=0}^{\infty} \int_{t_n}^{t_{n+1}} x^2(t) dt = \sum_{n=0}^{\infty} 1 = \boxed{\infty = E_x}$$

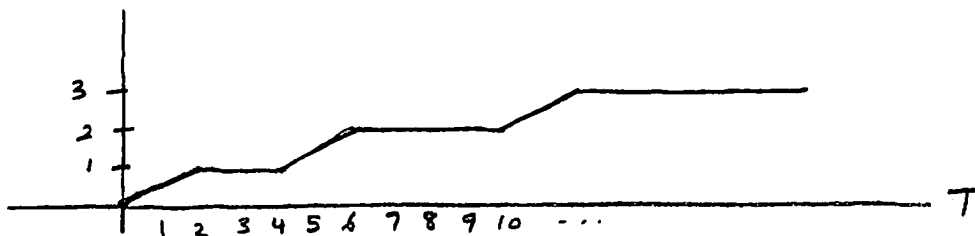
where  $t_n = 0, 2, 5, 9, 14, \dots$

Now power...

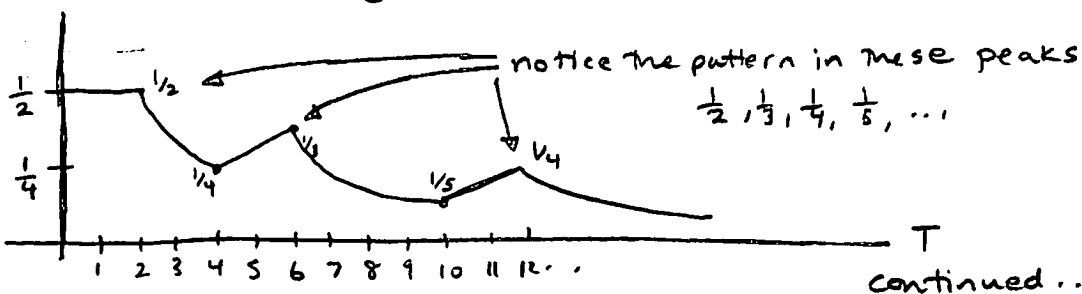
$$P_x = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} x^2(t) dt = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^{T/2} x^2(t) dt$$

because  $x(t) = 0$  for all  $t < 0$

Let's look at  $\int_0^{T/2} x^2(t) dt$



Now let's sketch  $\frac{1}{T} \int_0^{T/2} x^2(t) dt$



continued...

problem 1 continued...

The trend is evident here, but we can formalize our result by noting that

$$\frac{1}{T} \int_0^{T/2} x^2(t) dt \leq \frac{1}{n} \quad \text{for} \quad (n-1)! \leq T \leq n!$$

e.g. suppose  $n=2$ , we see that  $\frac{1}{T} \int_0^{T/2} x^2(t) dt \leq \frac{1}{2}$   
 for  $1 \leq T \leq 2$  from the graph.

Hence  $\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^{T/2} x^2(t) dt$  must be zero.

$$\Rightarrow P_x = 0.$$

2. a) Use KVL to write

$$x(t) = Ri(t) + y(t)$$

We know that  $i(t) = C \frac{dy(t)}{dt}$

$$\Rightarrow RC \frac{dy(t)}{dt} + y(t) = x(t)$$

in Lathi's notation  $(RC D + 1)y(t) = x(t)$

b) Following the approach on pp. 164-165,  
 we have  $Q(D) = RC D + 1$  and  $P(D) = 1$

We need to rewrite these polynomials in the correct form by dividing both sides by RC.

diff equation:  $(D + \frac{1}{RC})y(t) = \frac{1}{RC}x(t)$

$$Q(D) = D + \frac{1}{RC} \quad \text{with} \quad a_1 = \frac{1}{RC}$$

$$P(D) = 0 \cdot D + \frac{1}{RC} \quad \text{with} \quad b_0 = 0 \quad \text{and} \quad b_1 = \frac{1}{RC}$$

Since  $b_0 = 0$ ,  $h(t)$  consists only of characteristic modes.

The single root of  $Q(D)$  is  $\lambda + \frac{1}{RC} = 0 \Rightarrow \lambda = -\frac{1}{RC}$

Hence  $h(t) = b e^{-\frac{1}{RC}t} u(t)$ . We need to find  $b$ ...

To find  $K_1$ , we let  $x(t) = \delta(t)$  and  $y(t) = h(t)$  to write

$$\dot{h}(t) + \frac{1}{RC} h(t) = \frac{1}{RC} \delta(t)$$

initial condition:  $h(0^-)$  is zero (relaxed initial cond.)

but  $h(0^+) = K_1$ . Hence  $\dot{h}(0) = K_1 \delta(t)$ .

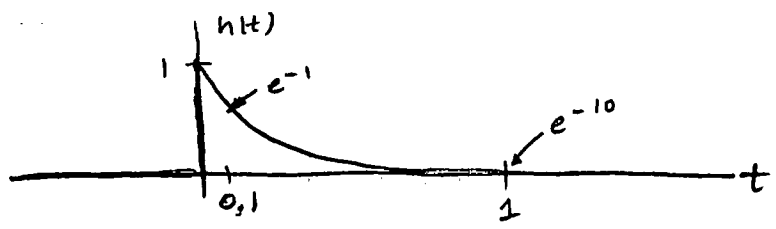
Match coefficients of impulse terms on both sides

$$K_1 \delta(t) + (\text{not an impulse}) = \frac{1}{RC} \delta(t)$$

Hence  $K_1 = \frac{1}{RC}$ .

Hence  $h(t) = \frac{1}{RC} e^{-\frac{1}{RC}t} u(t)$  is the impulse response.

When  $R=1$  and  $C=0.1$ ,  $h(t) = e^{-10t} u(t)$



c) Step response can be found by convolution

$$y(t) = \int_{-\infty}^{\infty} x(\tau) h(t-\tau) d\tau \quad x(\tau) = u(\tau)$$

↑  
impulse response from part b

The steps follow those in the 23-Mar lecture

$$y(t) = \int_0^{\infty} h(t-\tau) d\tau = \begin{cases} \int_0^t \frac{1}{RC} e^{-\frac{1}{RC}(t-\tau)} d\tau & t \geq 0 \\ 0 & \text{otherwise.} \end{cases}$$

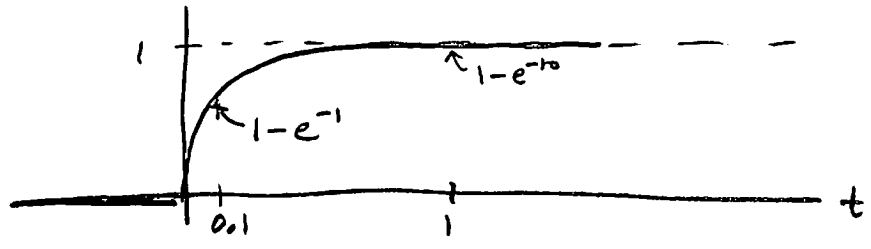
$$= \begin{cases} e^{-\frac{t}{RC}} \int_0^t \frac{1}{RC} e^{\frac{\tau}{RC}} d\tau & t \geq 0 \\ 0 & \text{otherwise} \end{cases} = \begin{cases} e^{-\frac{t}{RC}} \left[ e^{\frac{t}{RC}} - 1 \right] & t \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

continued...

$$y(t) = \begin{cases} 1 - e^{-t/Rc} & t \geq 0 \\ 0 & \text{otherwise.} \end{cases} = [1 - e^{-t/Rc}] u(t)$$

This is the step response.

When  $R=1$  and  $C=0.1$ , we have



3. Lathi 2.4-6

$$\int_{-\infty}^{\infty} u(\tau) \sin(t-\tau) u(t-\tau) d\tau = \sin t u(t) * u(t)$$

From  $u(\tau)$  and  $u(t-\tau)$  we have the limits of integration between 0 and  $t$

$$= \begin{cases} \int_0^t \sin(t-\tau) d\tau & t \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

$$= \begin{cases} 1 - \cos(t) & t \geq 0 \\ 0 & \text{otherwise.} \end{cases}$$

$$\cos t u(t) * u(t) = \int_{-\infty}^{\infty} u(\tau) \cos(t-\tau) u(t-\tau) d\tau$$

Same idea...

$$= \begin{cases} \int_0^t \cos(t-\tau) d\tau & t \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

$$= \begin{cases} \sin(t) & t \geq 0 \\ 0 & \text{otherwise.} \end{cases}$$

5. Lathi 2.4-7  $h(t) = e^{-t} u(t)$

a)  $x(t) = u(t)$

$$y(t) = (1 - e^{-t}) u(t)$$

This was done in lecture, but you can do this with line 2 of the table setting  $\lambda = -1$ .

b)  $x(t) = e^{-t} u(t)$  use line 5 of the table with  $\lambda = -1$

$$y(t) = t e^{\lambda t} u(t) = t e^{-t} u(t)$$

c)  $x(t) = e^{-2t} u(t)$  use line 4 of the table with  $\lambda_1 = -2$  and  $\lambda_2 = -1$

$$y(t) = \frac{e^{\lambda_1 t} - e^{\lambda_2 t}}{\lambda_1 - \lambda_2} u(t) = (e^{-t} - e^{-2t}) u(t)$$

d)  $x(t) = \sin(3t) u(t)$

Use Euler's identity  $\sin x = \frac{1}{2j} [e^{jx} - e^{-jx}]$

We can use line 2 and linearity to compute the result.

$$x_1(t) = \frac{1}{2j} e^{j3t} u(t) \rightarrow \boxed{S} \rightarrow y_1(t) = \frac{1}{2j} \left[ \frac{e^{j3t} - e^{-t}}{j3 + 1} \right] u(t)$$

$h(t) = e^{-t} u(t)$

$$x_2(t) = \frac{-1}{2j} e^{-j3t} u(t) \rightarrow \boxed{S} \rightarrow y_2(t) = \frac{-1}{2j} \left[ \frac{e^{-j3t} - e^{-t}}{-j3 + 1} \right] u(t)$$

Note that we've used the homogeneity property here. Now let's use additivity

$$x(t) = \sin(3t) u(t) = x_1(t) + x_2(t)$$

$$y(t) = y_1(t) + y_2(t) = \frac{1}{2j} \left[ \frac{e^{j3t} - e^{-t}}{j3 + 1} - \frac{e^{-j3t} - e^{-t}}{-j3 + 1} \right] u(t)$$

Simplify...

Note that  $(1+j3)(1-j3) = 1+9 = 10$

$$\text{So } y(t) = \frac{1}{20j} \left[ (e^{j3t} - e^{-t})(1-j3) - (e^{-j3t} - e^{-t})(1+j3) \right] u(t)$$

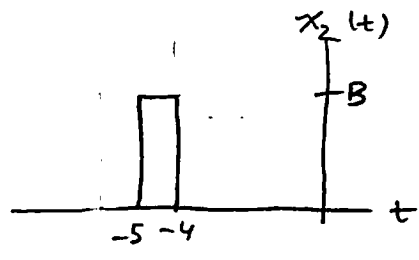
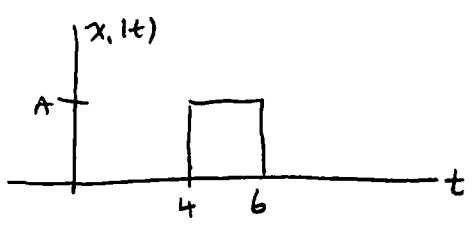
$$= \frac{1}{20j} \left[ e^{j3t} - e^{-t} - e^{-j3t} + e^{-t} + j3(-e^{j3t} + e^{-t} - e^{-j3t} + e^{-t}) \right] u(t)$$

$$= \frac{1}{20j} \left[ e^{j3t} - e^{-j3t} \right] u(t) + \frac{3}{20} \left[ 2e^{-t} - e^{-j3t} - e^{-j3t} \right] u(t)$$

$$y(t) = \left( \frac{1}{10} \sin(3t) + \frac{3}{10} e^{-t} - \frac{3}{10} \cos(3t) \right) u(t)$$

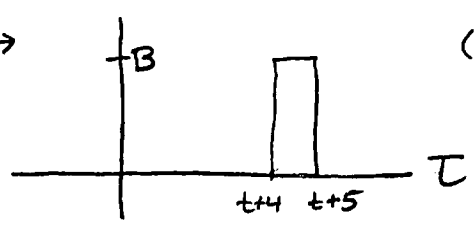
6. Lathi 2.4-18

a)



$$y(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$$

draw  $x_2(t-\tau)$  as a function of  $\tau$

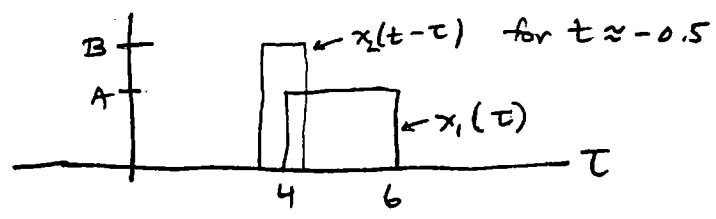


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it is clear that there is no overlap until  $t+5 \ge 4$  or, equivalently,  $t \ge -1$

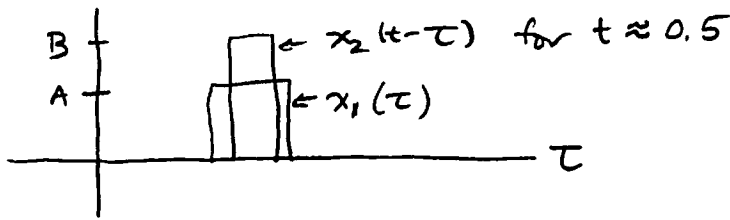
So, for  $t < -1$ ,  $y(t) = 0$

For  $-1 \le t \le 0$  The area under the product of  $x_1(\tau) x_2(t-\tau)$  is increasing linearly to a value of  $AB$  at  $t = 0$ .

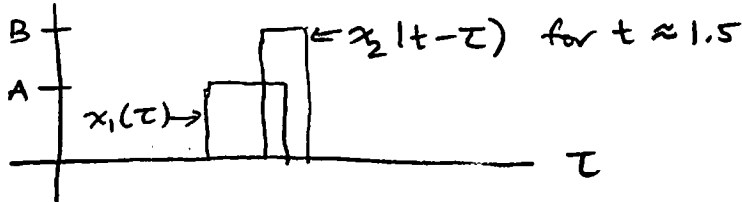


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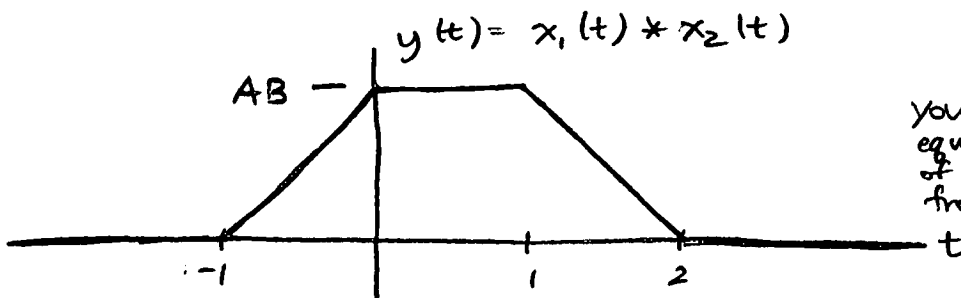
for  $0 \leq t \leq 1$  The area under the product of  $x_1(\tau) x_2(t-\tau)$  remains constant (equal to  $AB$ )



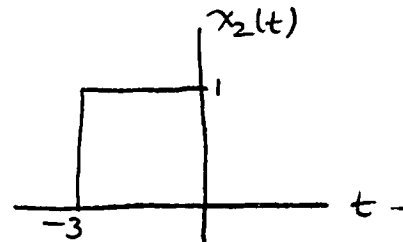
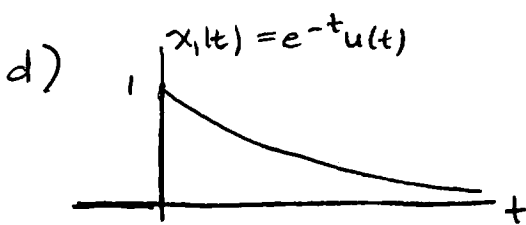
for  $0 \leq t \leq 2$ , the area under the product of  $x_1(\tau) x_2(t-\tau)$  decreases from  $AB$  to zero linearly



for  $t > 2$ , there is no overlap, hence the output is zero.  
put it all together to get

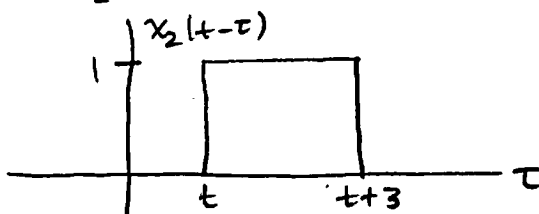


You can write equations for each of these lines directly from the graph



$$y(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$$

draw  $x_2(t-\tau)$  as a function of  $\tau$



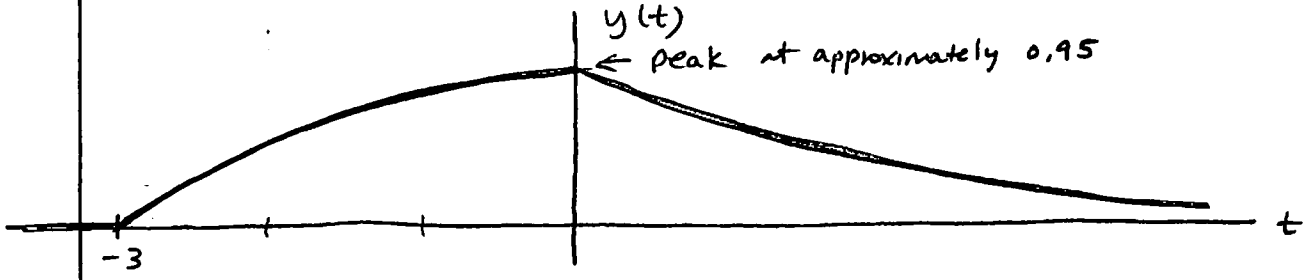
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continued...

it should be clear that there is no overlap until  $t+3 \geq 0$  or, equivalently,  $t \geq -3$

as  $t$  increases from  $-3$ , the area under the curve increases, but by a little less each time. When  $t$  becomes zero, the area under the product  $x_1(\tau)x_2(t-\tau)$  will peak and then decrease for all remaining time as  $t \rightarrow \infty$ .

So we can sketch the result



We can compute this output analytically using the fact that

$$x(t) = u(t+3) - u(t)$$

From the lecture, we know that

$$u(t) \rightarrow \boxed{e^{-t}u(t)} \rightarrow (1 - e^{-t})u(t) = y_1(t)$$

and time invariance implies

$$u(t+3) \rightarrow \boxed{e^{-t}u(t)} \rightarrow (1 - e^{-(t+3)})u(t+3) = y_2(t)$$

$$\text{so } y(t) = y_2(t) - y_1(t) = (1 - e^{-(t+3)})u(t+3) - (1 - e^{-t})u(t)$$

$$y(t) = \begin{cases} 0 & \text{when } t < -3 \\ 1 - e^{-(t+3)} & \text{when } -3 \leq t < 0 \\ e^{-t} - e^{-(t+3)} & \text{when } t \geq 0 \end{cases}$$

you can plot this in Matlab to confirm the sketch.

(end).