

# Digital Signal Processing

## IIR Filter Design via Impulse Invariance

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# Basic Procedure

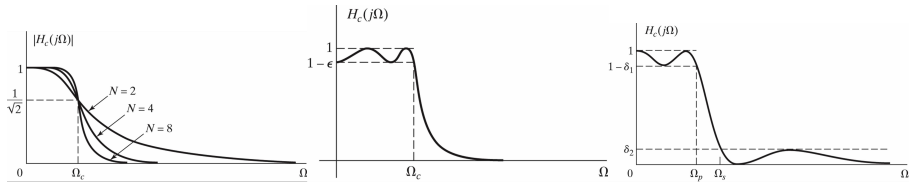
We assume here that we've already decided to use an IIR filter.

The basic procedure for IIR filter design via impulse invariance is:

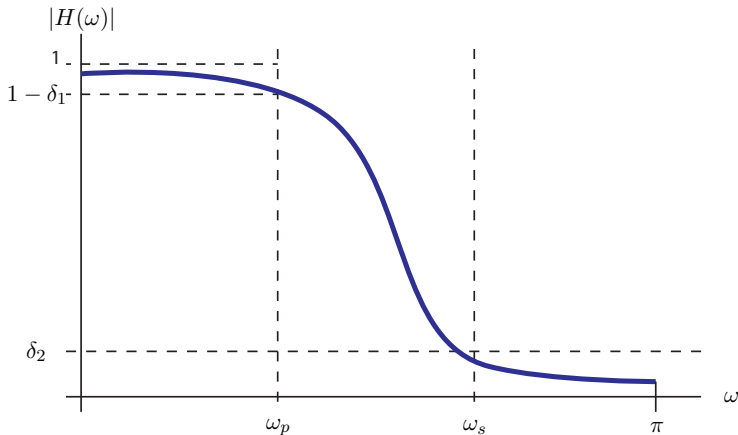
1. Determine the CT filter design method:
  - 1.1 Butterworth
  - 1.2 Chebychev Type I or Type II
  - 1.3 Elliptic
  - 1.4 ...
2. Transform the DT filter specifications to CT (sampling period  $T_d$  is arbitrary)
3. Design CT filter based on the magnitude squared response  $|H_c(j\Omega)|^2$ 
  - ▶ Determine filter order
  - ▶ Determine cutoff frequency
4. Determine  $H_c(s) \leftrightarrow h_c(t)$  corresponding to a stable causal filter
5. Convert to DT filter  $H(z) \leftrightarrow h[n]$  via impulse invariance such that  $h[n] = h_c(nT_d)$

# Determining the Continuous-Time Filter Design Method

|                   | passband   | stopband   |
|-------------------|------------|------------|
| Butterworth       | monotonic  | monotonic  |
| Chebyshev Type I  | equiripple | monotonic  |
| Chebyshev Type II | monotonic  | equiripple |
| elliptic          | equiripple | equiripple |



# Impulse-Invariant Lowpass Butterworth Filter Design Ex.



We start with the desired specifications of the DT filter. For this example, we will use  $\omega_p = 0.2\pi$ ,  $\omega_s = 0.3\pi$ ,  $1 - \delta_1 = 0.89125$ , and  $\delta_2 = 0.17783$ .

# Convert DT Filter Specs to CT Filter Specs

We will use the Butterworth filter approach in this example. A CT Butterworth filter has a squared magnitude response given by

$$|H_c(j\Omega)|^2 = \frac{1}{1 + \left(\frac{j\Omega}{j\Omega_c}\right)^{2N}} \quad (1)$$

where  $\Omega_c$  is the cutoff frequency (radians/second) and  $N$  is the filter order.

When designing filters via impulse invariance, the sampling period  $T_d$  is arbitrary. It is often convenient to just set  $T_d = 1$  so that  $\Omega = \omega$ .

This implies the CT filter specs can be written as  $\Omega_p = 0.2\pi$  and  $\Omega_s = 0.3\pi$ . Along with our magnitude specifications  $1 - \delta_1 = 0.89125$  and  $\delta_2 = 0.17783$ , we can substitute these results directly into (1) to write

$$\begin{aligned} 1 + \left(\frac{0.2\pi}{\Omega_c}\right)^{2N} &\leq \left(\frac{1}{0.89125}\right)^2 \\ 1 + \left(\frac{0.3\pi}{\Omega_c}\right)^{2N} &\geq \left(\frac{1}{0.17783}\right)^2 \end{aligned}$$

## Determine the $N$ and $\Omega_c$

We can take the previous inequalities and write them as equalities as

$$1 + \left( \frac{0.2\pi}{\Omega_c} \right)^{2N} = \left( \frac{1}{0.89125} \right)^2 \quad (2)$$

$$1 + \left( \frac{0.3\pi}{\Omega_c} \right)^{2N} = \left( \frac{1}{0.17783} \right)^2. \quad (3)$$

We have two equations and two unknowns. By taking logarithms, we can isolate  $N$  and  $\Omega_c$  to get

$$N = 5.8858$$

$$\Omega_c = 0.70474$$

Note that  $N$  must be an integer, so we can choose  $N = 6$ . We now can decide whether to pick  $\Omega_c$  to match the passband spec (and exceed the stopband spec) or match the stopband spec (and exceed the passband spec). To minimize the effect of aliasing, we usually choose the former.

## Determine $H_c(s)$ (part 1 of 3)

Given  $N = 6$  and our choice to match the passband spec, we have the equality

$$1 + \left( \frac{0.2\pi}{\Omega_c} \right)^{12} = \left( \frac{1}{0.89125} \right)^2 \quad (4)$$

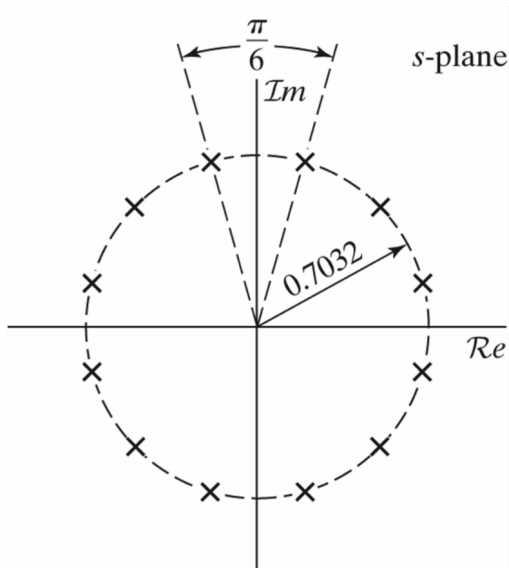
which gives  $\Omega_c = 0.7032$ . Now, to determine  $H_c(s)$ , we can write

$$H_c(s)H_c(-s) = \frac{1}{1 + \left( \frac{s}{j\Omega_c} \right)^{2N}}.$$

The pole locations of  $H_c(s)H_c(-s)$  follow from the fact that  $1 + \left( \frac{x}{a} \right)^M = 0$  has  $M$  solutions given by

$$x = \{ae^{j\pi/M}, ae^{j3\pi/M}, \dots, ae^{j(\pi+(M-1)2\pi)/M}\}.$$

The  $\frac{M}{2}$  poles corresponding to  $H_c(s)$  are those in the left half plane.

Determine  $H_c(s)$  (part 2 of 3)



# Determine $H_c(s)$ (part 3 of 3)

Choosing the poles from the left half plane and doing a little bit of algebra, we can write

$$H_c(s) = \frac{0.12093}{(s^2 + 0.3640s + 0.4945)(s^2 + 0.9945s + 0.4945)(s^2 + 1.3585s + 0.4945)}$$

Since all of the poles are simple (non-repeated), we can write the partial fraction expansion

$$H_c(s) = \sum_{k=1}^6 \frac{A_k}{s - s_k}$$

which implies that

$$h_c(t) = \begin{cases} \sum_{k=1}^6 A_k e^{s_k t} & t \geq 0 \\ 0 & t < 0. \end{cases}$$

# Determine $H(z)$ via Impulse Invariance

Impulse invariance requires  $h[n] = h_c(nT_d)$  where we have previously chosen  $T_d = 1$ . Hence we have

$$h[n] = \begin{cases} \sum_{k=1}^6 A_k e^{s_k n} & n \geq 0 \\ 0 & n < 0. \end{cases}$$

and it follows that

$$H(z) = \sum_{k=1}^6 \frac{A_k}{1 - e^{s_k} z^{-1}}$$

with ROC  $|z| > \text{largest magnitude pole}$ . A bit of algebra yields the final result

$$\begin{aligned} H(z) = & \frac{0.2871 - 0.4466z^{-1}}{1 - 1.2971z^{-1} + 0.6949z^{-2}} + \frac{-2.1428 + 1.1455z^{-1}}{1 - 1.0691z^{-1} + 0.3699z^{-2}} \\ & + \frac{1.8577 - 0.6303z^{-1}}{1 - 0.9972z^{-1} + 0.2570z^{-2}} \end{aligned}$$

which can immediately be realized in parallel form or rearranged to be realized in cascade or direct forms.

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